

# Smith Chart & Waveguides

Module V comprehensive notes – graphical analysis of transmission lines, impedance matching, waveguide modes and field solutions.

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## SECTION 01

### The Smith Chart

The Smith Chart is a graphical tool for analysing transmission lines. It gives a visual representation of how the impedance of a transmission line and its reflection coefficient change as we move along the line – without requiring complex calculations.

The chart consists of a circular plot with families of interlaced circles. Its main advantage is that transmission line characteristics can be determined without complex computational work.

#### KEY IDEA

##### Normalised Impedance

Since every TL has a different  $Z_0$ , we normalise all impedances:  $z_L = Z_L / Z_0$ .

This means a single Smith Chart works for any transmission line.

#### REFLECTION COEFFICIENT

##### $\Gamma_L$ Formula

The chart is built from  $\Gamma_L = (Z_L - Z_0) / (Z_L + Z_0)$ , which in normalised form becomes  $\Gamma_L = (z_L - 1) / (z_L + 1)$ .

#### BOUNDARY POINTS

### Short & Open Circuits

Left extreme:  $r=0, x=0 \rightarrow$  *short circuit*.

Right extreme:  $r=\infty, x=\infty \rightarrow$  *open circuit*.

## Reflection Coefficient on the Chart

Both  $z_L$  and  $\Gamma_L$  are complex. Writing  $z_L = r + jx$  and  $\Gamma_L = u + jv$ , we get:

#### DERIVATION

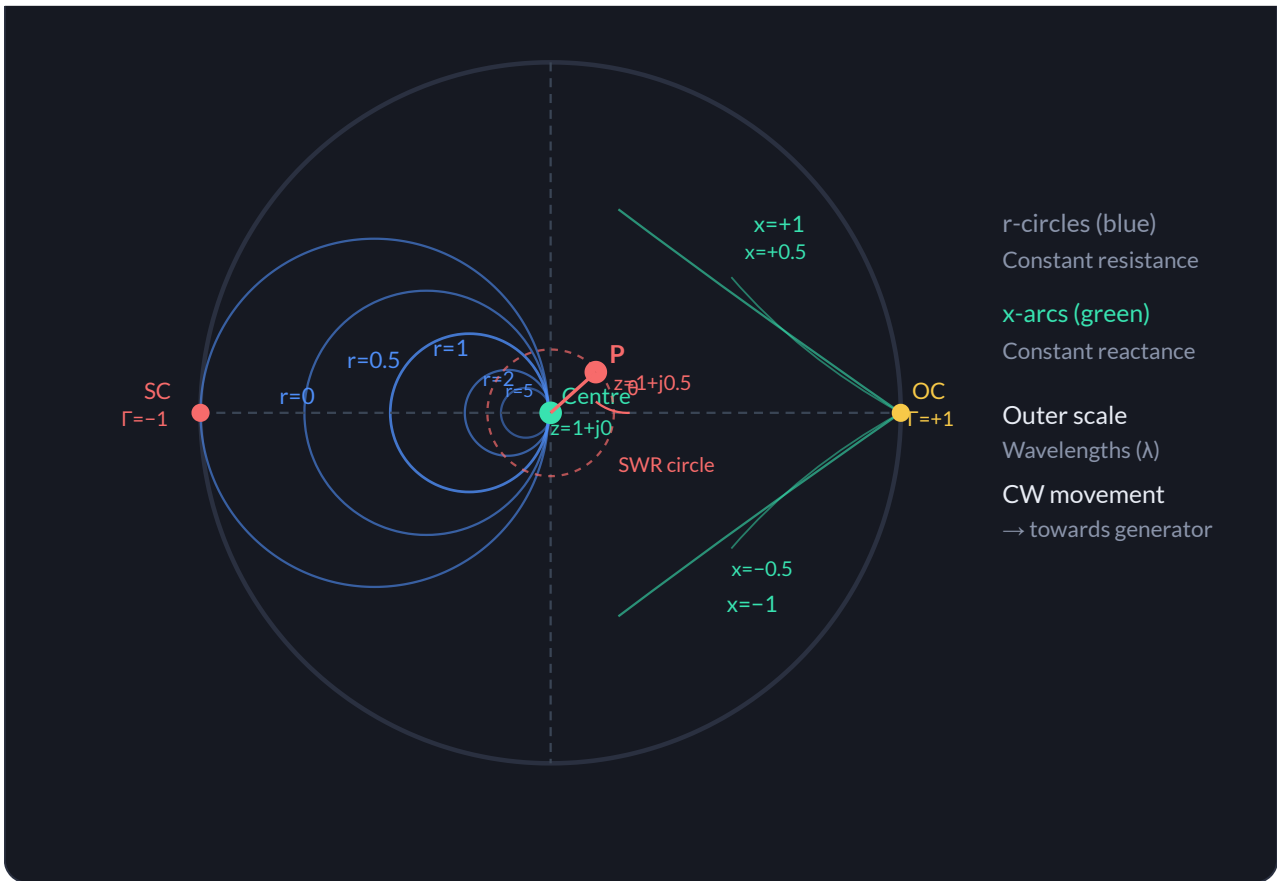
$$z_L = (1 + \Gamma_L) / (1 - \Gamma_L)$$

$$r = (1 - u^2 - v^2) / (1-u)^2 + v^2$$

$$x = 2v / (1-u)^2 + v^2$$

The Smith Chart lies within the unit circle:  $|\Gamma_L| \leq 1$ . The reflection coefficient is written as  $\Gamma_L = |\Gamma_L| e^{j\theta} = \Gamma_r + j\Gamma_i$ .

#### SMITH CHART – CONCEPTUAL LAYOUT



## Important Points about the Smith Chart

POINT 1

**Revolution =  $\lambda/2$**

A complete  $360^\circ$  revolution around the chart equals a distance of  $\lambda/2$  on the transmission line. Clockwise = towards generator. Counter-clockwise = towards load.

POINT 2

**Three Scales**

Outermost: distance from generator (in  $\lambda$ ). Middle: distance from load (in  $\lambda$ ). Innermost: angle protractor for measuring  $\theta$  of  $\Gamma$ .

POINT 3

**Dual Use**

The Smith Chart can serve as both an impedance chart and an admittance chart. The r and x circles correspond to g and b circles when used for admittance.

SECTION 02

# Construction — r and x Circles

## Resistance Circles (r-circles)

By manipulating the real part of  $z_L = (1+\Gamma)/(1-\Gamma)$ , we arrive at the equation of a circle:

R-CIRCLE EQUATION

$$(u - r/(1+r))^2 + v^2 = (1/(1+r))^2$$

Centre:  $(r/(1+r), 0)$  Radius:  $1/(1+r)$

NORMALISED R

RADIUS

CENTRE (U, V)

0

1

(0,0)

0.5	2/3	(1/3, 0)
1	1/2	(1/2, 0)
2	1/3	(2/3, 0)
5	1/6	(5/6, 0)
$\infty$	0	(1, 0)

All r-circles pass through the point (1,0) on the right of the chart. As  $r \rightarrow \infty$  the circles collapse to that single point.

## Reactance Circles (x-circles)

Similarly, from the imaginary part:

### X-CIRCLE EQUATION

$$(u - 1)^2 + (v - 1/x)^2 = (1/x)^2$$

Centre: (1, 1/x) Radius: |1/x|

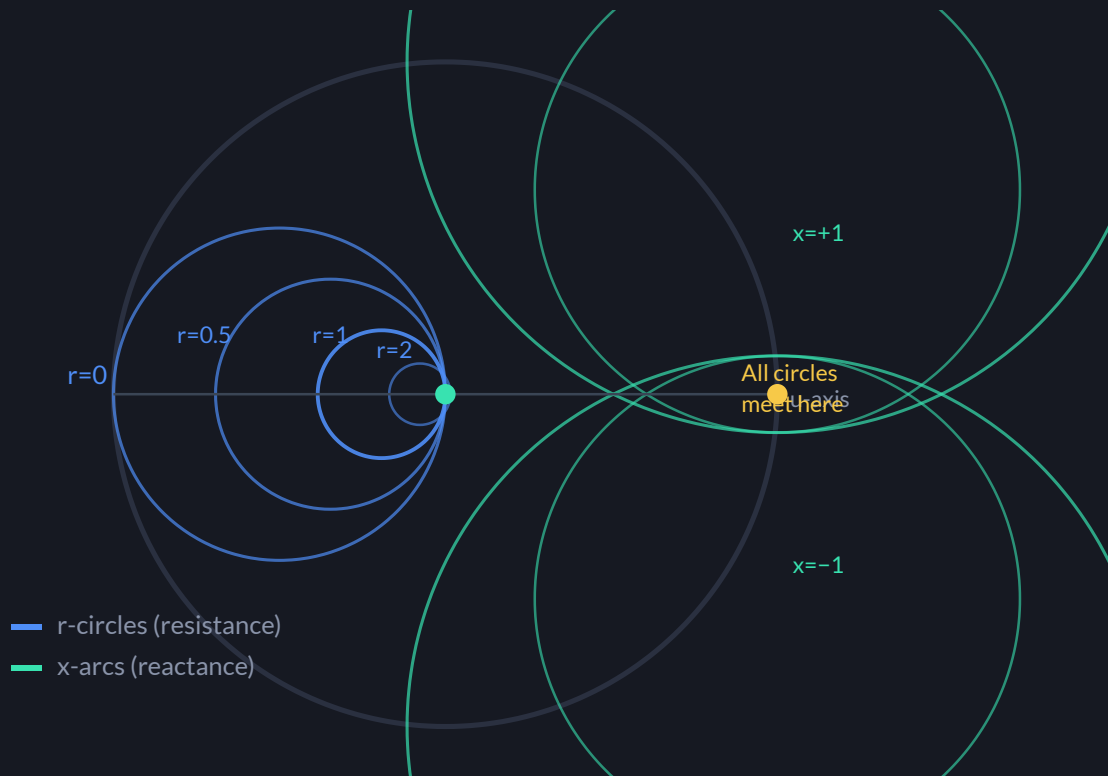
NORMALISED X	RADIUS	CENTRE
0	$\infty$	(1, $\infty$ ) – horizontal axis
$\pm 0.5$	2	(1, $\pm 2$ )
$\pm 1$	1	(1, $\pm 1$ )
$\pm 2$	1/2	(1, $\pm 1/2$ )
$\pm 5$	1/5	(1, $\pm 1/5$ )
$\pm \infty$	0	(1, 0)



All x-circles also pass through (1,0). Positive x (inductive) arcs are in the upper half; negative x (capacitive) arcs are in the lower half. Superimposing r-circles and x-circles

onto the unit circle gives the complete Smith Chart.

### R-CIRCLES AND X-ARCS – SUPERIMPOSED



## SECTION 03

# Applications of the Smith Chart

- Plotting complex impedance  $Z_L = R + jX$  on the chart
- Finding reflection coefficient  $\Gamma_L$  (magnitude and phase) for a given load
- Finding the Standing Wave Ratio (SWR) for any impedance
- Finding the admittance  $Y_L$  of a given impedance  $Z_L$
- Finding the input impedance  $Z_{in}$  at any point along a transmission line

→ Matching a transmission line to a load using a single series stub

### (a) Plotting Impedance

Normalise the load impedance:  $z_L = Z_L/Z_0$ . Locate the intersection of the r-circle and x-arc corresponding to the real and imaginary parts.

**i** Example:  $Z_L = 300 - j25 \Omega$  on a  $50 \Omega$  line  $\rightarrow z_L = 6 - j0.5$ . Since  $x = -0.5$ , the point Q is in the lower half of the chart at the intersection of the  $r=6$  circle and  $x=-0.5$  arc.

### (b) Finding Reflection Coefficient $\Gamma_L$

#### PROCEDURE

1. Normalise and locate point P on chart
2. Draw line from centre O to P
3. Extend to meet  $r=0$  circle at point G
4.  $|\Gamma_L| = OP / OG$  (OG corresponds to  $|\Gamma|=1$ )
5.  $\theta =$  angle of OP with horizontal axis

Example:  $Z_L=60+j40$ ,  $Z_0=50 \rightarrow z_L=1.2+j0.8 \rightarrow |\Gamma|=0.35\angle 56^\circ$

### (c) Standing Wave Ratio (SWR)

Once  $z_L$  is located (say at point P), draw a circle centred at O with radius OP. This is the constant SWR circle (also called the constant  $|\Gamma|$  circle).

#### SWR READING

Locate point S where the SWR circle meets the positive real axis ( $\Gamma_r$ -axis).

The value of  $r$  at S = SWR value

If  $Z_L=Z_0$ ,  $z_L=1+j0 \rightarrow |\Gamma|=0 \rightarrow \text{SWR}=1$  (matched line)

### (d) Load Admittance

If  $z_L$  is marked at point P, extend line OP through the centre to the diametrically opposite point P'. The value read at P' gives the normalised admittance  $y_L$ .

#### EXAMPLE (FROM NOTES)

$$Z_L = 60 + j40, Z_0 = 50 \Omega$$

$$z_L = 1.2 + j0.8 \rightarrow P \text{ on chart}$$

$$y_L \text{ at } P' = 0.575 - j0.38$$

$$Y_L = Y_0 \cdot y_L = (1/50)(0.575 - j0.38) = 0.0115 - j0.0076 \text{ S}$$

### (e) Input Impedance $Z_{in}$

To find  $Z_{in}$  at a distance  $l$  from the load:

#### STEP-BY-STEP

1. Normalise:  $z_L = Z_L/Z_0$ , locate P
2. Convert  $l$  to degrees:  $360^\circ \leftrightarrow \lambda/2$ , so  $l(\text{in } \lambda) \times 720^\circ = \text{angle}$
3. Rotate clockwise by that angle on SWR circle  $\rightarrow$  new point R
4. Read  $z_{in}$  at R, then  $Z_{in} = Z_0 \cdot z_{in}$

Example:  $l=0.4\lambda$  on  $Z_0=75\Omega$  line,  $Z_L=100+j150\Omega$

$$z_L = 1.33 + j2 \rightarrow \text{rotate } 288^\circ \text{ CW} \rightarrow z_{in} = 0.29 + j0.64$$

$$Z_{in} = 75(0.29 + j0.64) = (21.75 + j48) \Omega$$

## SECTION 04

# Transmission Line Sections as Circuit Elements

From the input impedance formula  $Z_{in} = Z_0 [(Z_L + jZ_0 \tan \beta l) / (Z_0 + jZ_L \tan \beta l)]$ , special cases arise for open and short-circuited lines.

## Open-Circuited Line ( $Z_L = \infty$ )

RESULT

$$Z_{in} = -jZ_0 \cot(\beta l)$$

Pure reactance. Capacitive when  $\beta l < \pi/2$ ; inductive when  $\pi/2 < \beta l < \pi$ .

## Short-Circuited Line ( $Z_L = 0$ )

RESULT

$$Z_{in} = +jZ_0 \tan(\beta l)$$

Pure reactance. Capacitive when  $\pi/2 < \beta l < \pi$ ; inductive when  $\beta l < \pi/2$ .



Both open- and short-circuited TL sections can act as inductor or capacitor depending on their length. No physical L or C components are needed at microwave frequencies!

## Half-Wavelength Section ( $\lambda/2$ )

$$L = \lambda/2 \rightarrow \beta L = \pi \rightarrow \tan(\pi) = 0$$

$$Z_{in} = Z_0 \cdot [Z_L / Z_0] = Z_L$$

A half-wavelength section is transparent – the input impedance equals the load impedance. Useful when physical spacing between two ports is needed without altering electrical behaviour.

## Quarter-Wavelength Section ( $\lambda/4$ ) – Transformer

$$L = \lambda/4 \rightarrow \beta L = \pi/2 \rightarrow \tan(\pi/2) = \infty$$

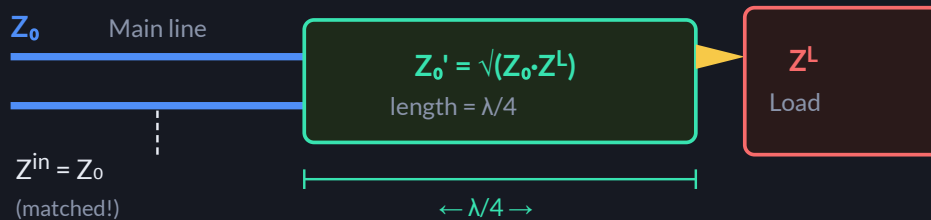
$$Z_{in} = Z_0^2 / Z_L$$

$$\therefore Z_0 = \sqrt{(Z_{in} \cdot Z_L)}$$

The quarter-wave transformer (also called quarter-wave matching section) is used for impedance matching. To match a load  $Z_L$  to a line  $Z_0$ , insert a  $\lambda/4$  section with characteristic impedance  $Z_0' = \sqrt{(Z_0 \cdot Z_L)}$ .

! Example: To match a  $120 \Omega$  load to a  $75 \Omega$  line, insert a  $\lambda/4$  section with  $Z_0' = \sqrt{(75 \times 120)} = 95 \Omega$ .

#### QUARTER-WAVE TRANSFORMER – IMPEDANCE MATCHING



#### SECTION 05

## Waveguides

A waveguide is a hollow conducting tube used to transmit EM waves from one point to another. Unlike transmission lines, it can support many field configurations (modes).

### TL vs Waveguide Comparison

PROPERTY

TRANSMISSION LINE

WAVEGUIDE

Supported modes	TEM wave only	TE, TM (many modes)
Frequency range	DC to very high f	Only above cutoff $f_c$
Efficiency at microwave	Low (skin effect, dielectric loss)	High (less loss)
Bandwidth	Narrower	Larger
Filter behaviour	—	High-pass (cannot transmit DC)

! Waveguides act as high-pass filters — they can only transmit frequencies above a minimum cutoff frequency. Below cutoff, the wave is attenuated and does not propagate.

## SECTION 06

# Modes of Propagation

The waves propagating through a waveguide have infinite possible patterns of electric and magnetic fields — these patterns are called **modes**. Three fundamental mode types exist:

### MODE 1

#### TEM — Transverse Electromagnetic

Both  $E_z = 0$  and  $H_z = 0$ . Both E and H are perpendicular to the direction of propagation. Cannot exist in a hollow waveguide — requires two conductors.

### MODE 2

#### TE — Transverse Electric (H-wave)

$E_z = 0$ ,  $H_z \neq 0$ . E is entirely transverse; H has a longitudinal component. Called  $TE_{mn}$  mode ( $m, n =$  mode indices).

### MODE 3

#### TM — Transverse Magnetic (E-wave)

$H_z = 0$ ,  $E_z \neq 0$ . H is entirely transverse; E has a longitudinal component. Called

TM<sub>mn</sub> mode.

#### MODE SUMMARY

TEM:  $E_z = 0, H_z = 0$

TE:  $E_z = 0, H_z \neq 0$  (dominant in rectangular WG)

TM:  $H_z = 0, E_z \neq 0$

#### SECTION 07

## Parallel Plate Waveguide

Two perfectly conducting parallel plates separated by distance 'a', of infinite extent in y and z directions. The medium between them is lossless dielectric ( $\sigma=0$ ) with parameters  $\epsilon$  and  $\mu$ .

### Boundary Conditions

BC 1

#### Tangential $E = 0$

$E_t = 0$  at the conducting surface. Electric field terminates normally on the conductor.

BC 2

#### Normal $H = 0$

$H_n = 0$  at the conducting surface. Magnetic field is tangential to the conductor.

### General Field Equations (from Maxwell)

WITH ASSUMPTIONS: PROPAGATION IN Z, UNIFORM IN Y ( $\partial/\partial Y = 0$ )

$$E_x = -(\gamma/h^2) \cdot \partial E_z / \partial x \quad [A]$$

$$H_x = -(\gamma/h^2) \cdot \partial H_z / \partial x \quad [B]$$

$$E_y = (j\omega\mu/h^2) \cdot \partial H_z / \partial x \quad [C]$$

$$H_y = (-j\omega\epsilon/h^2) \cdot \partial E_z / \partial x \quad [D]$$

$$\text{where } h^2 = \gamma^2 + \mu\epsilon\omega^2$$

## TE Wave Solution ( $E_z=0$ )

Solving with BCs  $E_y=0$  at  $x=0$  and  $x=a$  gives sinusoidal variation:

### TE FIELD COMPONENTS (PARALLEL PLATES)

$$E_x = 0 \quad E_z = 0$$

$$E_y = A_1 \sin(m\pi x/a) e^{-j\beta z}$$

$$H_x = -(\beta/j\mu\omega) A_1 \sin(m\pi x/a) e^{-j\beta z}$$

$$H_z = -(A_1 m\pi/j\mu\omega a) \cos(m\pi x/a) e^{-j\beta z}$$

$$h = m\pi/a, \text{ where } m = \pm 1, \pm 2, \pm 3 \dots (m \neq 0)$$

Lowest order mode:  $TE_{10}$

## TM Wave Solution ( $H_z=0$ )

### TM FIELD COMPONENTS (PARALLEL PLATES)

$$E_y = 0 \quad H_x = 0$$

$$E_x = (\gamma A_4/j\epsilon\omega) \cos(m\pi x/a) e^{-\gamma z}$$

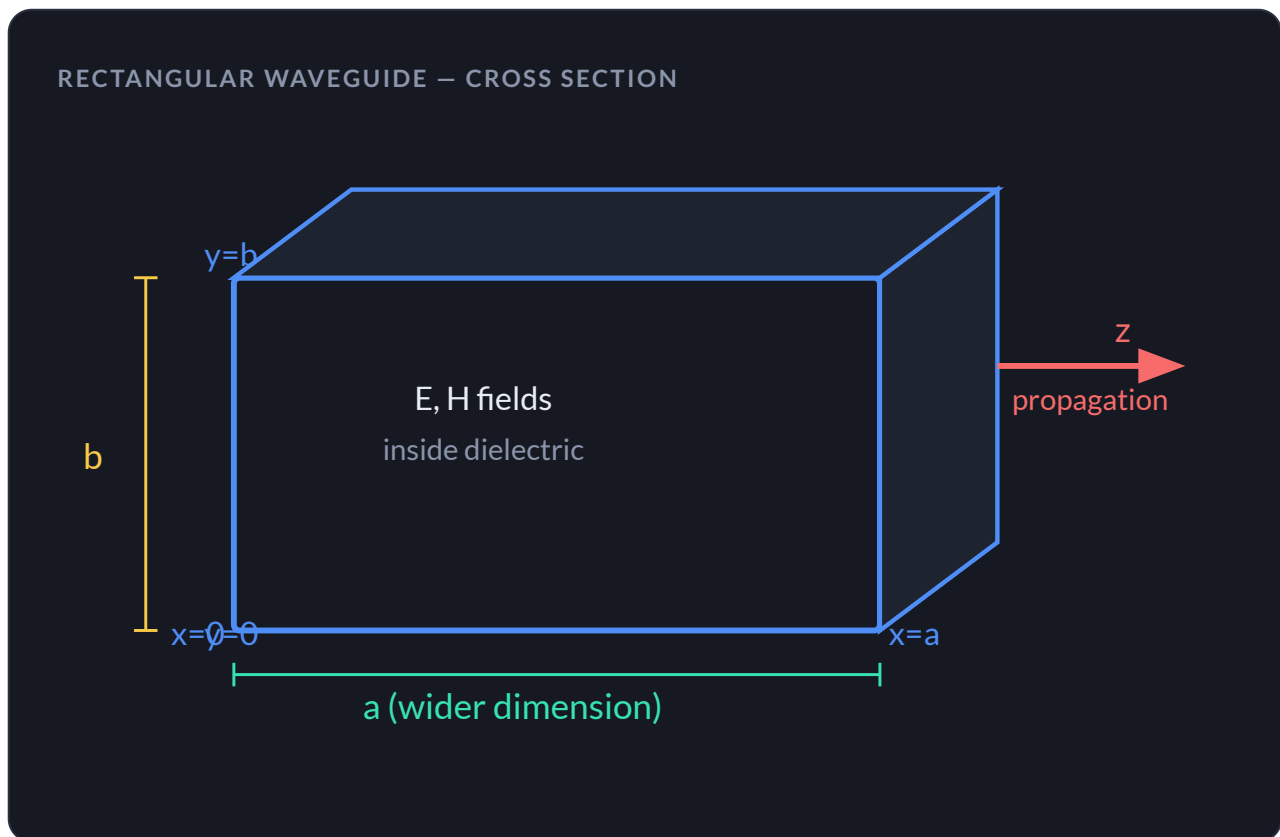
$$E_z = -(h A_4/j\epsilon\omega) \sin(m\pi x/a) e^{-\gamma z}$$

$$H_y = A_4 \cos(m\pi x/a) e^{-\gamma z}$$

Lowest order TM mode:  $TM_0$  ( $m=0$  allowed for TM)

# Rectangular Waveguide

A hollow rectangular conducting pipe with cross-section dimensions  $a$  (width) and  $b$  (height), propagating in the  $z$ -direction. More practical than infinite parallel plates.



## TE<sub>mn</sub> Mode in Rectangular Waveguide

For TE modes ( $E_z=0$ ), solving the wave equation by variable separation gives  $H_z(x,y,z)$ :

### H<sub>z</sub> SOLUTION

$$H_z(x,y,z) = A_{mn} \cos(m\pi x/a) \cos(n\pi y/b) e^{-j\beta z}$$

$m = 0, 1, 2, \dots$  and  $n = 0, 1, 2, \dots$ , not both zero

### ALL TE FIELD COMPONENTS

$$E_x = (j\omega\mu n\pi/h^2 b) A_{mn} \cos(m\pi x/a) \sin(n\pi y/b) e^{-j\beta z}$$

$$E_y = -(j\omega\mu m\pi/h^2 a) A_{mn} \sin(m\pi x/a) \cos(n\pi y/b) e^{-j\beta z}$$

$$H_x = (j\beta m\pi/h^2 a) A_{mn} \sin(m\pi x/a) \cos(n\pi y/b) e^{-j\beta z}$$

$$H_y = (j\beta n\pi/h^2 b) A_{mn} \cos(m\pi x/a) \sin(n\pi y/b) e^{-j\beta z}$$

## Cutoff Frequency

### CUTOFF FREQUENCY $F_c$

$$f_c = (v_p/2ab) \sqrt{[(mb)^2 + (na)^2]}$$

$$= (v_p/2) \sqrt{[(m/a)^2 + (n/b)^2]}$$

For air-filled WG:  $v_p = c = 3 \times 10^8$  m/s

$$\text{Cutoff wavelength: } \lambda_c = 2 / \sqrt{[(m/a)^2 + (n/b)^2]}$$

## Dominant Mode: $TE_{10}$

The  $TE_{10}$  mode ( $m=1, n=0$ ) has the lowest cutoff frequency and is the dominant mode in all rectangular waveguides. For  $TE_{10}$ :

### $TE_{10}$ CUTOFF FREQUENCY

$$f_c = v_p/(2a) \quad \lambda_c = 2a$$

## Phase and Group Velocities

### PHASE VELOCITY

$$v_p = \omega/\beta$$

Velocity of propagation of equiphase surfaces along the guide. Always greater than free-space velocity in a waveguide.

### GROUP VELOCITY

$$v_g = d\omega/d\beta$$

Velocity of energy propagation. Always less than free-space velocity. For lossless media:  $v_p \cdot v_g = v_c^2$

### KEY RELATIONSHIP

$$v_p \cdot v_g = v_c^2 = c^2$$

At cutoff:  $v_p \rightarrow \infty$ ,  $v_g \rightarrow 0$  (no energy propagation)

## Worked Example – Cutoff Frequencies

**Ex** Air-filled rectangular WG with  $a=100\text{mm}$ ,  $b=50\text{mm}$ . Find  $f_c$  for  $TE_{10}$ ,  $TE_{20}$ ,  $TE_{01}$ ,  $TE_{11}$ .

MODE	M	N	$F_c$
$TE_{10}$	1	0	1.5 GHz (dominant)
$TE_{20}$	2	0	3 GHz
$TE_{01}$	0	1	3 GHz
$TE_{11}$	1	1	3.354 GHz
$TE_{12}$	1	2	6.185 GHz
$TE_{21}$	2	1	4.243 GHz
$TM_{11}$	1	1	3.354 GHz (same as $TE_{11}$ )

### MODE CUTOFF FREQUENCY SPECTRUM (A=100MM, B=50MM)





$TE_{10}$  is the dominant mode because it has the lowest cutoff frequency. A waveguide is designed to operate in the  $TE_{10}$  mode by choosing the operating frequency to lie between  $f_c(TE_{10})$  and  $f_c(TE_{20})$ . This ensures only the dominant mode propagates — single-mode operation.

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